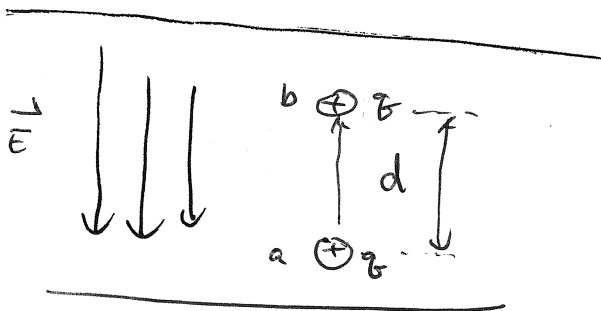


# ELECTRICAL POTENTIAL ENERGY

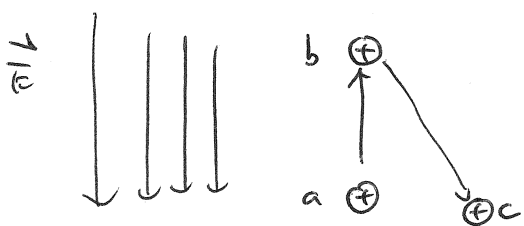


Work Done BY ELECTRIC FIELD

$$W_{a \rightarrow b} = -qEd$$

WORK DONE BY MY HAND

$$W = qEd$$



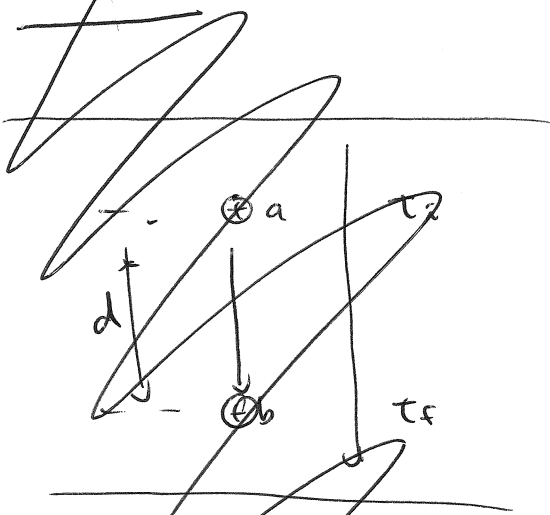
WORK DONE

$$W = qEd - qEd = 0$$

WORK ONLY DEPENDS ONLY ON FINAL POSITIONS

ELECTRIC FIELD  $\rightarrow$  CONSERVATIVE FORCES EXERTS  $\rightarrow$

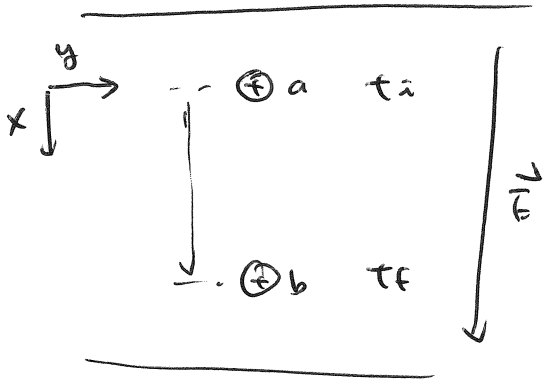
EXAMPLE



$$V(t_1) = 0$$

$$V(t_F) = ?$$

Now LET'S REVERSE THINGS A LITTLE



$$V(t_2) = 0 \quad \text{AT REST}$$

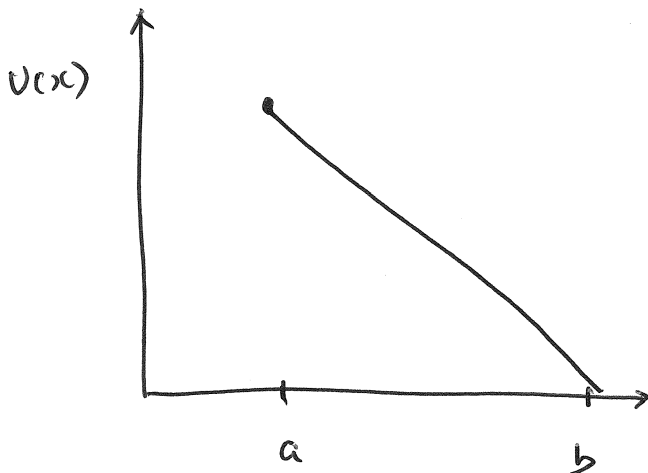
$$\vec{F} = q\vec{E}$$

$$W_{a \rightarrow b} = \text{WORK DONE BY THE FIELD BETWEEN } a \rightarrow b$$

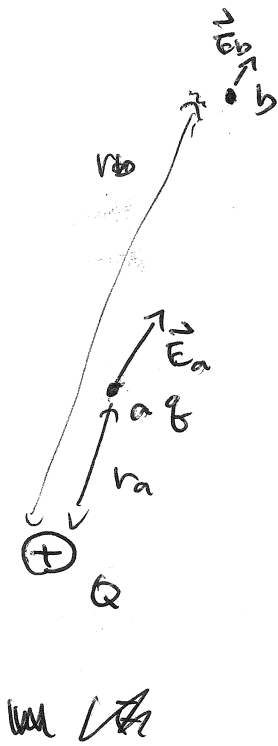
$$= qEd$$

IN CONSERVATIVE FIELD

$$W_{a \rightarrow b} = qEd = U_a - U_b = \text{POTENTIAL ENERGY}$$



ELECTRIC POTENTIAL ENERGY



WORK DONE BY THE FIELD

$$\int_{r_a}^{r_b} \vec{F} \cdot d\vec{r}$$

$$= \int_{r_a}^{r_b} \frac{1}{4\pi\epsilon_0} \frac{Qq}{r^2} dr$$

$$= \frac{Qq}{4\pi\epsilon_0} \left[ -\frac{1}{r} \right]_{r_a}^{r_b}$$

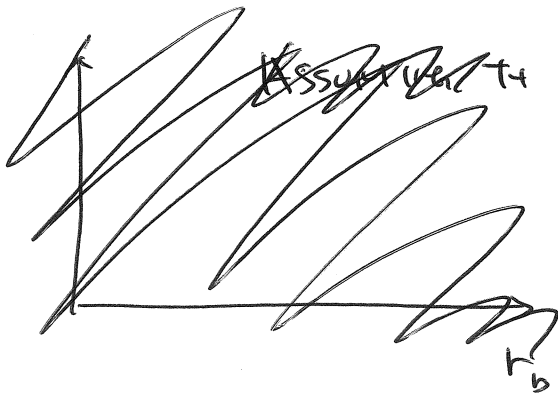
$$= \frac{Qq}{4\pi\epsilon_0} \left[ -\frac{1}{r_b} - \left( -\frac{1}{r_a} \right) \right]$$

$$W_{a \rightarrow b} = \frac{Qq}{4\pi\epsilon_0} \left[ \frac{1}{r_a} - \frac{1}{r_b} \right]$$

BY THE FIELD

$$W_{a \rightarrow b} = \frac{Qq_0}{4\pi\epsilon_0} \left[ \frac{1}{r_a} - \frac{1}{r_b} \right] = U_a - U_b$$

~~$\therefore V_{\text{ery}} = \frac{Qq_0}{4\pi\epsilon_0 r}$~~

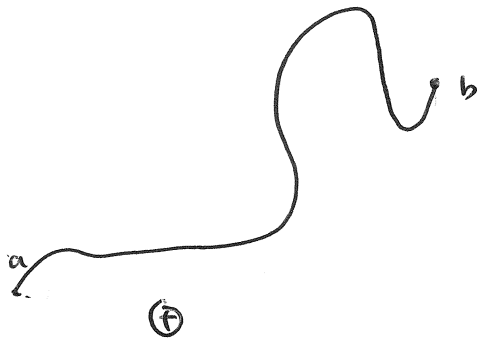


TAKE LIMIT  $r_b \rightarrow \infty$

$$U(r_b) = 0$$

ENERGY  
0 POTENTIAL, REALLY REALLY  
FAR AWAY

$$U(r) = \frac{Qq_0}{4\pi\epsilon_0 r} \quad \text{[Joules]}$$



$$W_{a \rightarrow b} = \frac{Qq_0}{4\pi\epsilon_0} \left[ \frac{1}{r_a} - \frac{1}{r_b} \right]$$

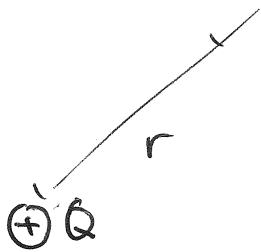
# ELECTRIC POTENTIAL

$$V = \frac{U}{q}$$

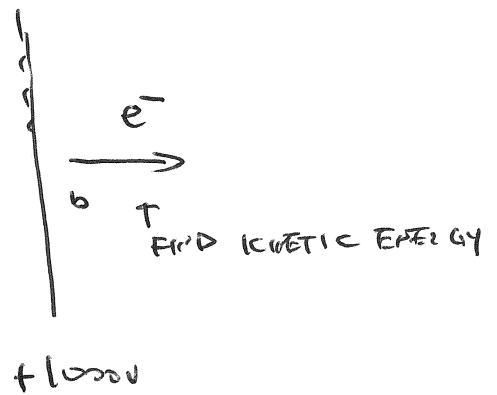
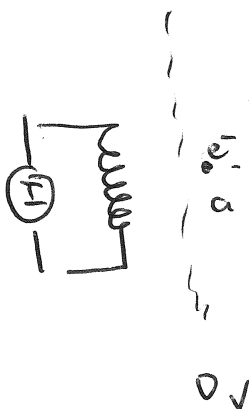
$$U = qV$$

POTENTIAL ENERGY PER UNIT CHARGE

V: UNITS VOLT



$$V(r) = \frac{Q}{4\pi\epsilon_0} \frac{1}{r} \quad [\text{VOLT}]$$



$$\frac{U(a) - U(b)}{q} = q \left[ \frac{V(a) - V(b)}{1} \right] = W_{a \rightarrow b} = k \cdot E \cdot (h)$$

$$\frac{V(a) - V(b)}{1.6 \times 10^{-19} \text{ C}} = \frac{100 \text{ V} - 0 \text{ V}}{1.6 \times 10^{-19} \text{ C}}$$

$k \cdot E = \frac{100 \text{ V}}{1.6}$

$$q [V(a) - V(b)] = -1.6 \times 10^{-19} \text{ C} \cdot -1000 \text{ V}$$

$$= 1.6 \times 10^{-16} \text{ V} \cdot \text{C}$$

$$= 1.6 \times 10^{-16} \text{ J}$$

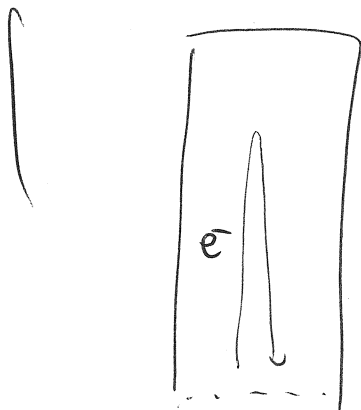
OR

$$1000 \text{ eV}$$

$$1 \text{ eV} = 1.6 \times 10^{-19} \text{ J}$$

---

ASIDE

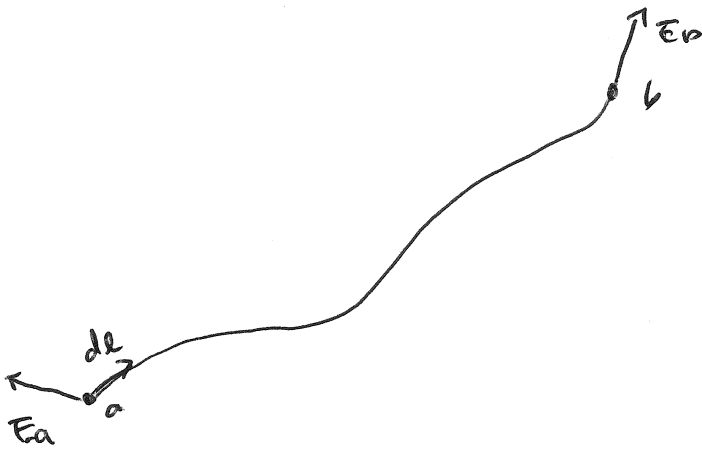


MASS OF ELECTRON IN ACCELERATED  
FRAME VS. INERTIAL FRAME

# REVAMP OUR UNDERSTANDING

WET

FRI



⊕ Q

q : "TEST" CHARGE

$W_{a \rightarrow b}$

(WORK DONE BY ELECTRIC FIELD ON CHARGE q)

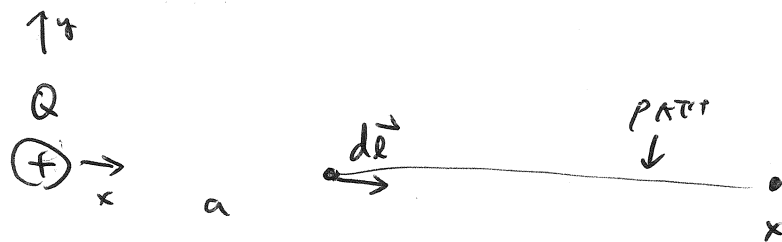
$$= \int_a^b \underset{\text{PATH}}{q \vec{E}} \cdot d\vec{\ell} = U(a) - U(b) = \underline{[V(a) - V(b)]} q$$

$$V(a) - V(b) = \int_a^b \vec{E} \cdot d\vec{\ell} \quad \text{(PATH)}$$

$$V(b) - V(a) = - \int_a^b \vec{E} \cdot d\vec{\ell} \quad \text{(PATH)}$$



SOUNDS HARD BUT CONSIDER



$$d\vec{l} = dr \hat{r}$$

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \hat{r}$$

$$V(x) - V(a) = - \int_a^x \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \hat{r} \cdot dr \hat{r}$$

$$= - \frac{1}{4\pi\epsilon_0} Q \int_a^x \frac{dr}{r^2} \quad \text{with } \hat{r} \cdot \hat{r} = 1$$

$$= - \left[ \frac{Q}{4\pi\epsilon_0} \left( -\frac{1}{r} \right) \right]_a^x$$

$$= \frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{x} - \frac{1}{a} \right] < 0$$

OK.

$$\frac{dV}{dx} = - \frac{Q}{4\pi\epsilon_0} \frac{1}{x^2} = -E_x(x)$$

$$\vec{E} = -\vec{\nabla}V = - \left[ \frac{\partial V}{\partial x} \hat{i} + \frac{\partial V}{\partial y} \hat{j} + \frac{\partial V}{\partial z} \hat{k} \right]$$

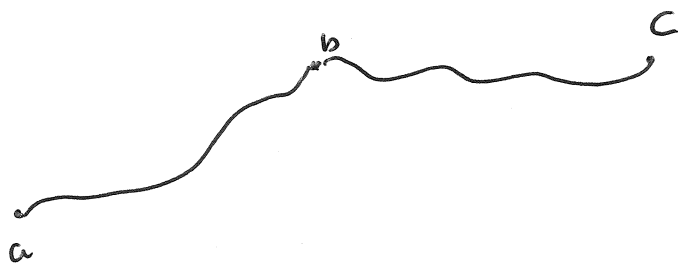
# NOTE

1.  $V$  : POTENTIAL NOT POTENTIAL ENERGY

2.  $\vec{E} = -\vec{\nabla}V$

$$\vec{\nabla} \times \vec{E} = -\vec{\nabla} \times \vec{\nabla}V = 0$$

3.



$$V(c) - V(a) = - \int_a^c \vec{E} \cdot d\vec{\ell}$$

$$= - \int_a^b \vec{E} \cdot d\vec{\ell} - \int_b^c \vec{E} \cdot d\vec{\ell}$$

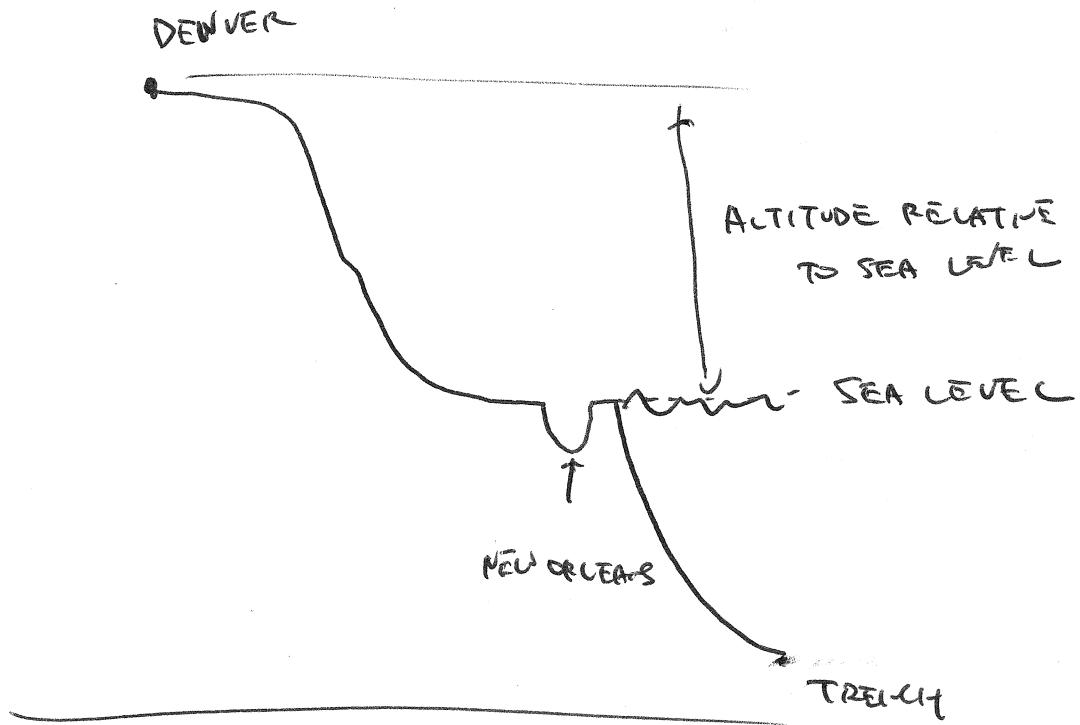
$$V(c) - V(a) = V(b) - V(a) - [V(c) - V(b)]$$

i.e.

$$(V_{ca} + \text{CONSTANT} = V_{bc})$$

IT'S ALL RELATIVE  
TO WHERE YOU START

MUCH LIKE ALTITUDE



TO STANDARDIZE WE SET  $V=0$  AT  $\infty$

$$V(\infty) = 0$$

↳ LAW OF SUPERPOSITION STILL APPLIES

↳ TEST CHANGE

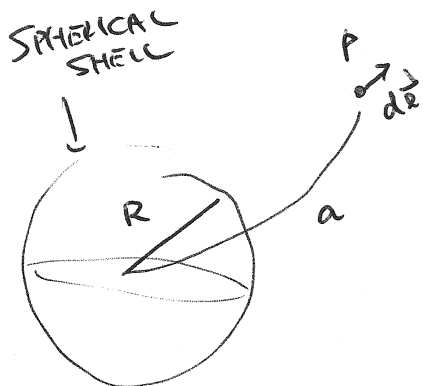
$\delta_1$   
 $\delta_2$   
 $\delta_3$

$$\vec{F}_T = \vec{F}_1 + \vec{F}_2 + \vec{F}_3$$

$$\vec{E} = \vec{E}_1 + \vec{E}_2 + \vec{E}_3$$

$$V = V_1 + V_2 + V_3$$

## EXAMPLE



OUTSIDE

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$$

INSIDE : NO CHARGE IS ENCLOSED

$$\vec{E} = 0$$

POTENTIAL :

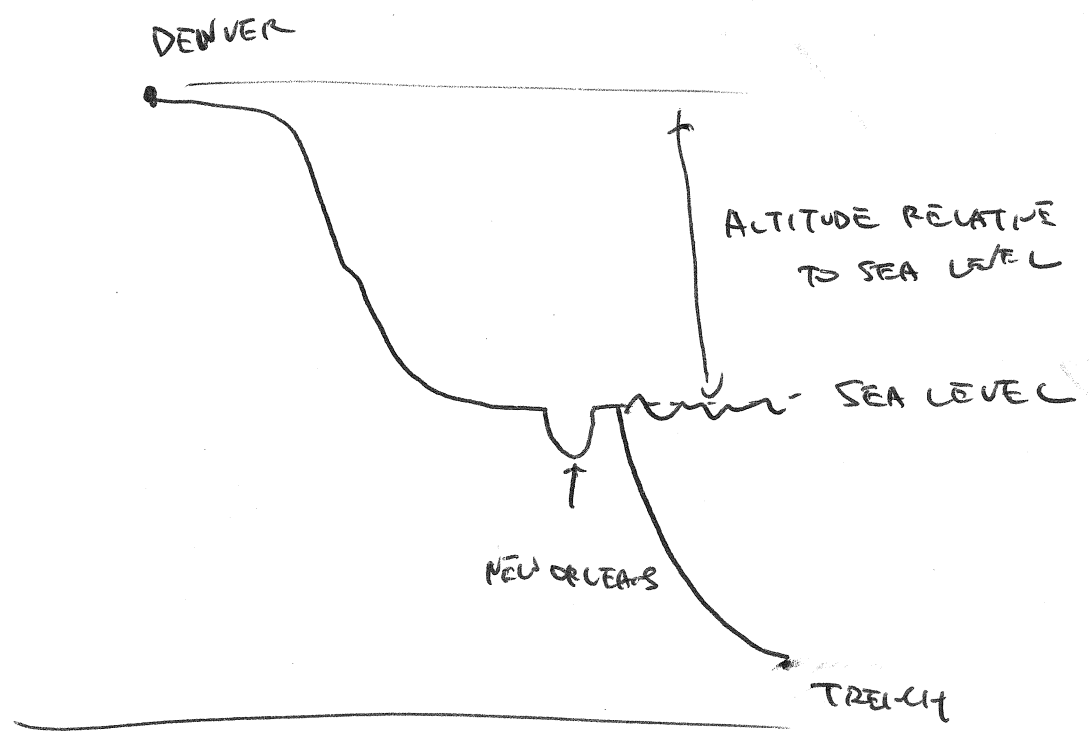
$$\begin{aligned} V(\infty) - V(a) &= 0 - V(a) = - \int_a^{\infty} \vec{E} \cdot d\vec{r} \\ &= - \int_a^{\infty} \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} dr \hat{r} \cdot \hat{r} \end{aligned}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{q}{r} \Big|_a^{\infty}$$

$$-V(a) = \frac{1}{4\pi\epsilon_0} q \left[ \frac{1}{\infty} - \frac{1}{a} \right]$$

$$V(a) = \frac{q}{4\pi\epsilon_0} \frac{1}{a} \quad (\text{For } a > R)$$

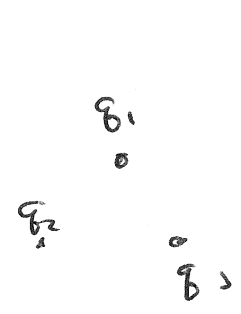
MUCH LIKE ALTITUDE



TO STANDARDIZE WE SET  $V=0$  AT  $\infty$

$$V(\infty) = 0$$

← LAW OF SUPERPOSITION STILL APPLIES



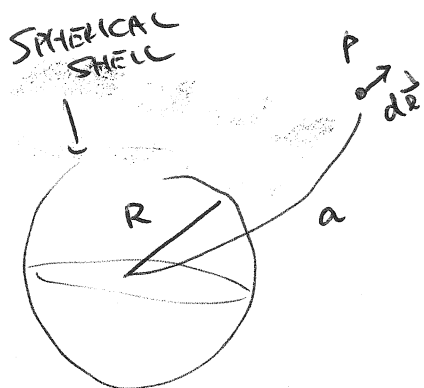
8 TEST CHARGE  
0

$$\vec{F}_T = \vec{F}_1 + \vec{F}_2 + \vec{F}_3$$

$$\vec{E} = \vec{E}_1 + \vec{E}_2 + \vec{E}_3$$

$$V = V_1 + V_2 + V_3$$

## EXAMPLE



OUTSIDE

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$$

INSIDE : NO CHARGE IS ENCLOSED

$$\vec{E} = 0$$

POTENTIAL :

$$\begin{aligned} V(\infty) - V(a) &= 0 - V(a) = - \int_a^{\infty} \vec{E} \cdot d\vec{r} \\ &= - \int_a^{\infty} \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} dr \hat{r} \cdot \hat{r} \\ &= \left[ \frac{1}{4\pi\epsilon_0} \frac{q}{r} \right]_a^{\infty} \end{aligned}$$

$$-V(a) = \frac{1}{4\pi\epsilon_0} q \left[ \frac{1}{\infty} - \frac{1}{a} \right]$$

$$V(a) = \frac{q}{4\pi\epsilon_0} \frac{1}{a} \quad (\text{For } a > R)$$

INSIDE :  $a < R$

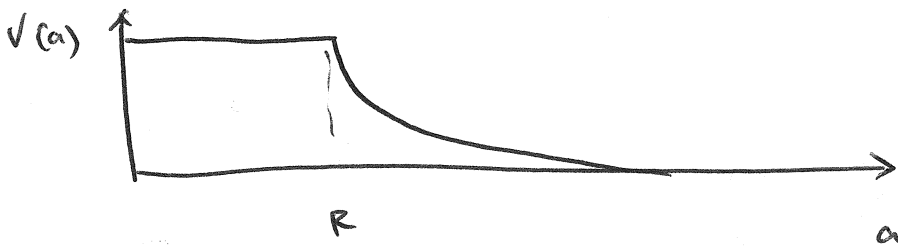
$$V(\infty) - V(a) = V(\infty) - V(R) + V(R) - V(a)$$

$$= 0 - \frac{1}{4\pi\epsilon_0} \frac{q}{R} + \int_a^R \vec{E} \cdot d\vec{\ell}$$

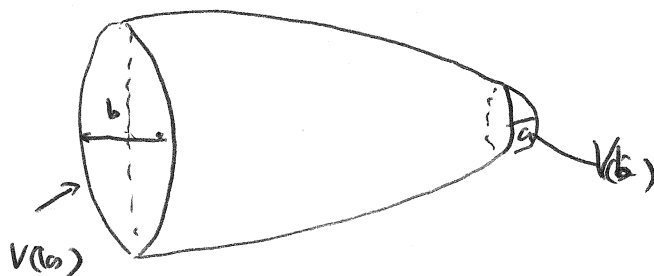
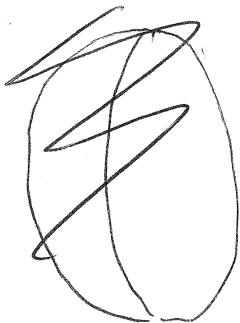
$E=0$  INSIDE

$$-V(a) = -\frac{1}{4\pi\epsilon_0} \frac{q}{R}$$

$$V(a) = \frac{1}{4\pi\epsilon_0} \frac{q}{R} \quad (\text{For } a < R)$$



ODD SHAPED OBJECT (METALLIC SHELL)



IF  $V(a) \neq V(b)$   $\Rightarrow$  CHARGE CAN MOVE AROUND

---

SO AT EQUILIBRIUM

$$V(a) = V(b)$$

$$\frac{q_a}{a} = \frac{q_b}{b}$$

$$\frac{q_a}{q_b} = \frac{a}{b}$$

$$\sigma_a \approx \frac{q_a}{a^2} \quad \sigma_b \approx \frac{q_b}{b^2}$$

$$\frac{a^2 \sigma_a}{b^2 \sigma_b} \approx \frac{a}{b}$$

$$\frac{\sigma_a}{\sigma_b} = \frac{b}{a}$$

$$b \gg a$$

$$\underline{\sigma_a \gg \sigma_b}$$

CHARGE IS CONCENTRATED

AT THE POINT

OF LARGE ~~CURVATURE~~

CURVATURE)



CORONA DISCHARGE : AIR CAN SUPPORT ONLY UP TO

$$3 \times 10^6 \text{ V/m}$$

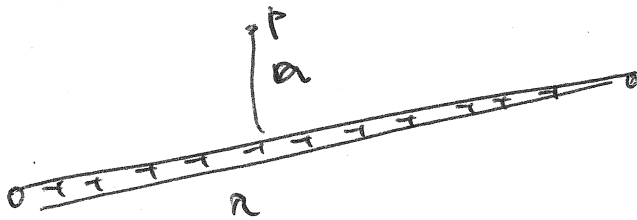


$$\vec{E}_{\text{SURFACE}} = \frac{V_{\text{SURFACE}}}{R}$$

$$3 \times 10^6 \text{ V/m} \cdot 0.1 = \underline{3 \times 10^5 \text{ V}}$$

DEMO REQUEST DISCHARGE

EXAMPLE #2



$$E(r) = \frac{1}{2\pi\epsilon_0} \frac{\lambda}{r} \hat{r} \quad (\text{GAUSS'S LAW}) \quad d\vec{x} = dr \hat{r}$$

$$V(r) - V(a) = - \int_a^r \vec{E} \cdot d\vec{x}$$

$$= - \int_a^r \frac{1}{2\pi\epsilon_0} \frac{\lambda}{r} \hat{r} \cdot dr \hat{r}$$

$$= - \int_a^r \frac{1}{2\pi\epsilon_0} \frac{\lambda}{r} dr$$

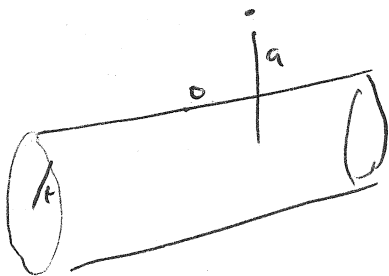
$$0 - V(a) = - \int_a^{\infty} \frac{\lambda}{2\pi\epsilon_0} \frac{1}{r} dr$$

$$= - \frac{\lambda}{2\pi\epsilon_0} \ln r \Big|_a^{\infty}$$

$$V(a) = \frac{\lambda}{2\pi\epsilon_0} [\infty - \ln a]$$

A PROBLEM

BUT WE CAN SET  $V=0$  AT ANY POINT



~~At R~~

$$V(a) - V(R) = - \int_R^a \frac{\lambda}{2\pi\epsilon_0} \frac{1}{r} dr$$

$$V(a) - 0 = - \frac{\lambda}{2\pi\epsilon_0} \ln r \Big|_R^a$$

$$= - \frac{\lambda}{2\pi\epsilon_0} (\ln a - \ln R)$$

$$V(a) = - \frac{\lambda}{2\pi\epsilon_0} \ln \frac{a}{R} = \frac{\lambda}{2\pi\epsilon_0} \ln \left( \frac{R}{a} \right)$$